

SIMPLE MODELS FOR SCALE DEPENDENT SPECTRAL DIMENSION

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SIMPLE MODELS FOR SCALE DEPENDENT SPECTRAL DIMENSION

1. Characterizing large scale structure of graphs
2. A curious result from quantum gravity
3. A definition
4. An easy example
5. An example using random graph ensembles
6. Conclusions

1. CHARACTERIZING LARGE SCALE STRUCTURE

Hausdorff dimension d_H -- we assume ∞ graphs

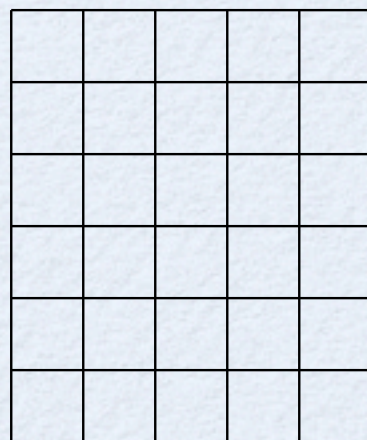
1. Choose a point r_0

2. Find all points $B_R(r_0)$ within graph distance R of r_0

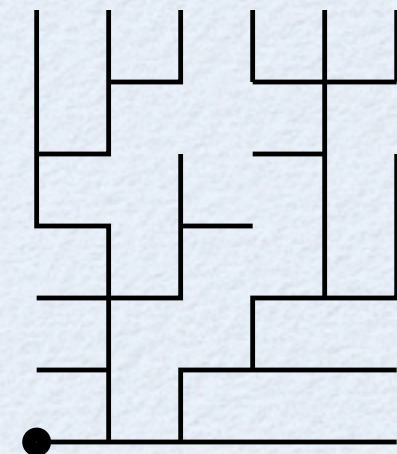
3. $|B_R(r_0)| \sim R^{d_H}$ as $R \rightarrow \infty$, independent of r_0

d_H is blind to some sorts of connectivity eg in the Euclidean metric

$d_H = 2$ for Z^2

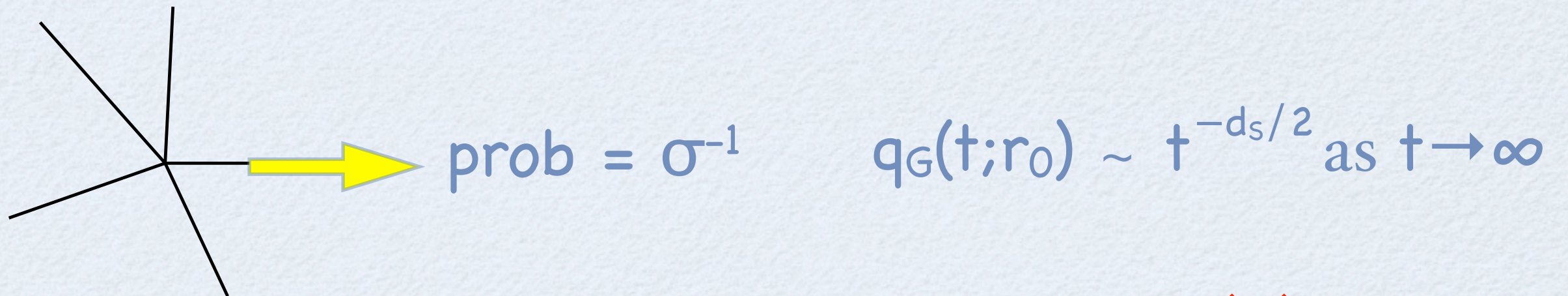


& Spanning Tree



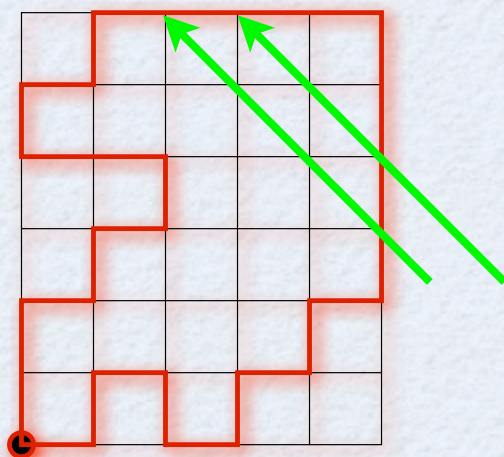
Spectral dimension d_S

1. Choose a point r_0
2. Random walker leaves r_0 at time 0 and returns at time t with probability $q_G(t; r_0)$

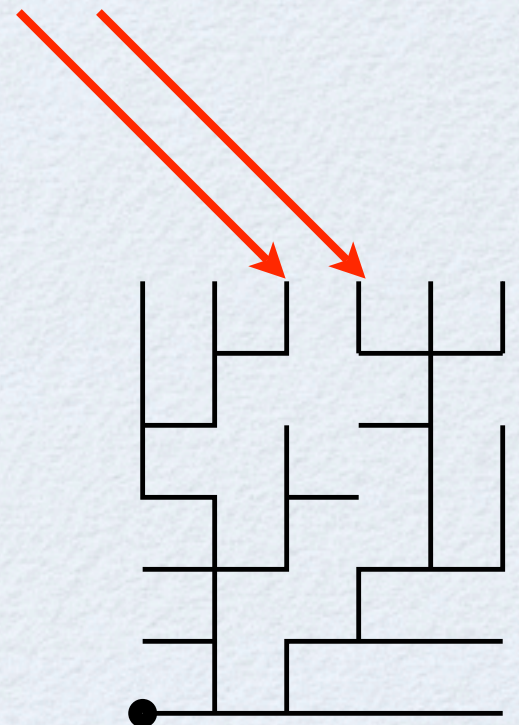


Random walk sees connectivity:

$d_S = 2$ for Z^2



but $16/13$ for



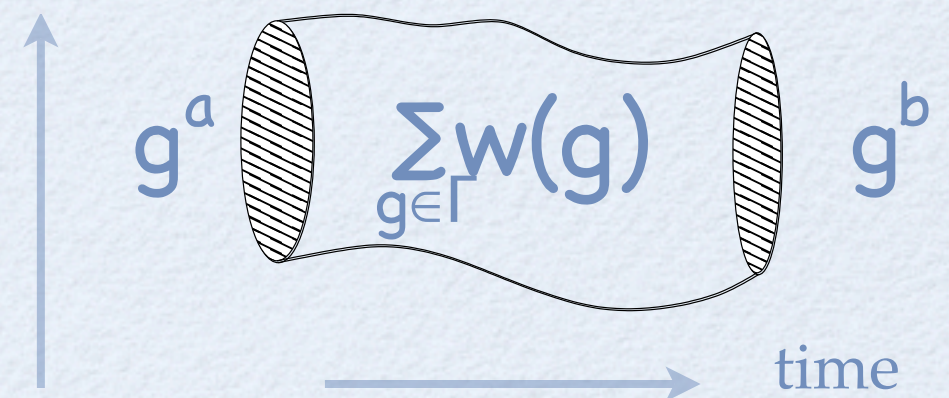
2. A CURIOUS RESULT FROM QUANTUM GRAVITY

Gravity's dynamical degree of freedom is the metric $g_{\mu\nu}(x,t)$

Classically $g_{\mu\nu}(x,t)$ obeys Einstein's equations:

$$g_{\mu\nu}(x,0) \longrightarrow g_{\mu\nu}(x,t)$$

Quantum mechanics is different:

$$\langle g^b(x), t=T \mid g^a(x), t=0 \rangle \sim \int_{g \in \Gamma} w(g)$$


Probability amplitude for evolution from g^a to g^b

How are Γ and w defined ?

Several approaches; some non-stringy ones

1. Continuum field theoretical -- exact RG looking for non-Gaussian fixed points where QG non-perturbatively renormalizable (asymptotic safety)

hep-th/0508202

2. Discretized -- Causal Dynamical Triangulations Γ is a set of graphs and we look for a critical point (or line) where a continuum limit can be taken to recover QG

hep-th/0505113

Curiosity is that both indicate $d_s=4$ at large distance scales
but $d_s=2$ at small distance scales

3. DEFINING SCALE DEPENDENT d_s

Convenient to use generating fn

$$Q_G(x, L) = 1 + \sum_{t=2}^{\infty} q_G(t; r_0) (1-x)^{t/2}$$
$$\sim x^{-1+d_s/2} \text{ as } x \rightarrow 0$$

Now define (if the limit exists),

$$\tilde{Q}(\xi, \lambda) = \lim_{a \rightarrow 0} a^{1/2} Q(a\xi, (\lambda/a)^\Delta)$$

And split the sum $\sum_{t=2}^{\infty} \dots = \sum_{t=2}^T \dots + \sum_{t=T}^{\infty} \dots$

Make suitable choices of T and bound one of the sums

Choose $T = \frac{1}{a\xi \log(1 + (\xi\lambda)^{-1})}$

$$\tilde{Q}(\xi, \lambda) (1 - e^{-\xi\lambda}) < \lim_{a \rightarrow 0} a^{1/2} \sum_{t=2}^T q_G(t) (1 - a\xi)^{t/2} < \tilde{Q}(\xi, \lambda)$$

$\tilde{Q}(\xi \rightarrow \infty, \lambda)$ describes walks of continuum duration $< \lambda$

Choose $T = \frac{\log(1 + \xi\lambda)}{a\xi}$

$$\tilde{Q}(\xi, \lambda) - ea^{1/2}\sqrt{T} < \lim_{a \rightarrow 0} a^{1/2} \sum_{t=T}^{\infty} q_G(t) (1 - a\xi)^{t/2} < \tilde{Q}(\xi, \lambda)$$

$\tilde{Q}(\xi \rightarrow 0, \lambda)$ describes walks of continuum duration $> \lambda$

So define spectral dimension at long distances by

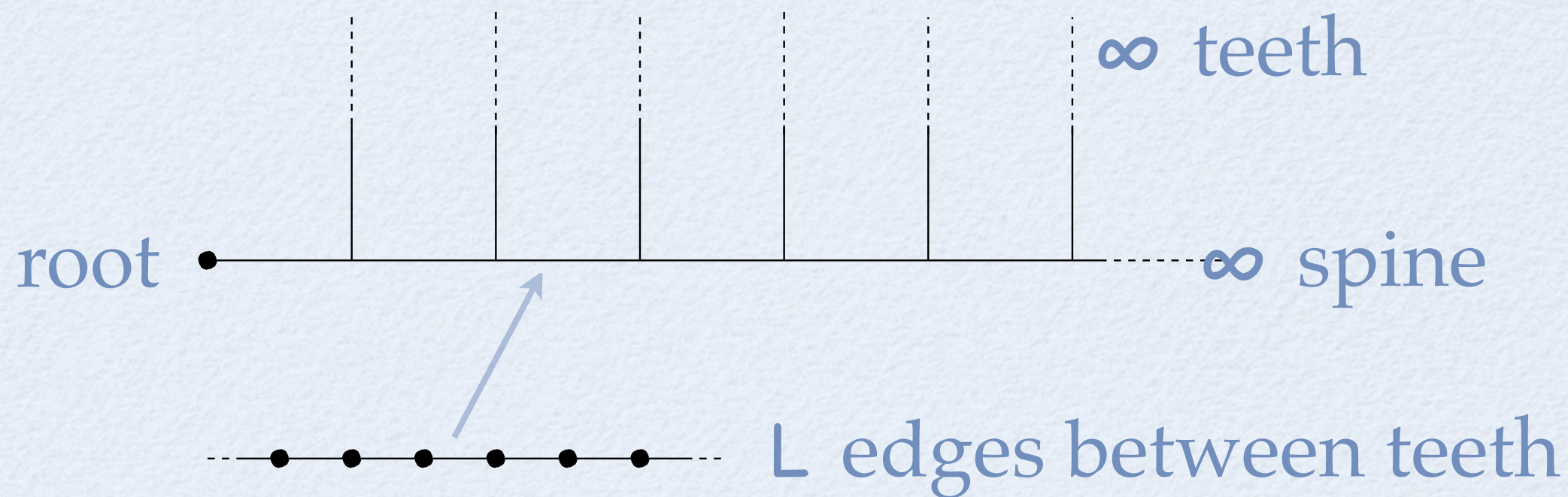
$$1 + d_s/2 = \lim_{\xi \rightarrow 0} \frac{\log \tilde{Q}(\xi, \lambda)}{\log \xi}$$

and at short distances by

$$1 + d_s/2 = \lim_{\xi \rightarrow \infty} \frac{\log \tilde{Q}(\xi, \lambda)}{\log \xi}$$

The existence and ordering of the limits are crucial

4. AN EASY EXAMPLE



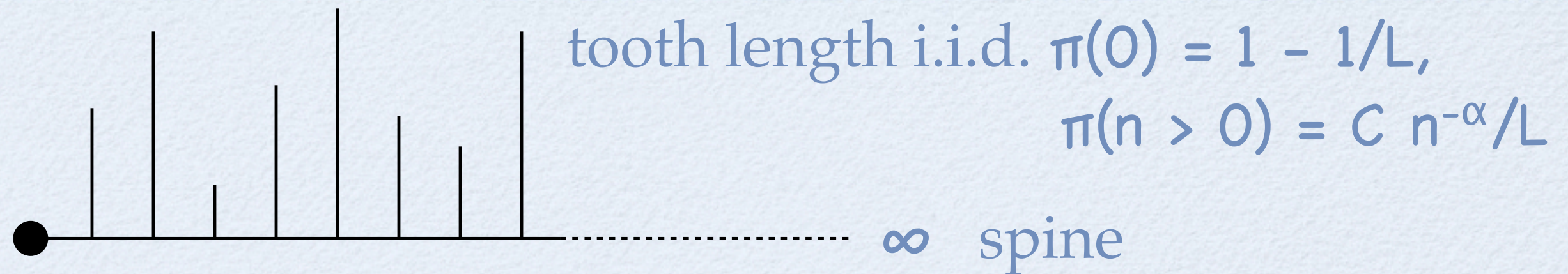
$Q(x,L)$ can be calculated exactly

$$\begin{aligned} \tilde{Q}(\xi,\lambda) &= \lim_{a \rightarrow 0} a^{1/2} Q(a\xi, (\lambda/a)^{1/2}) \\ &= \frac{2}{\xi^{1/2} (5 + 4 \coth(\lambda\xi)^{1/2})^{1/2}} \end{aligned}$$

$$\xi \rightarrow 0, d_s = 3/2$$

$$\xi \rightarrow \infty, d_s = 1$$

5. RANDOM ENSEMBLES after hep-th/0509191



Interesting range is $1 < \alpha \leq 2$; for L fixed

$$d_H = 3 - \alpha,$$

$$d_S = 2 - \alpha/2$$

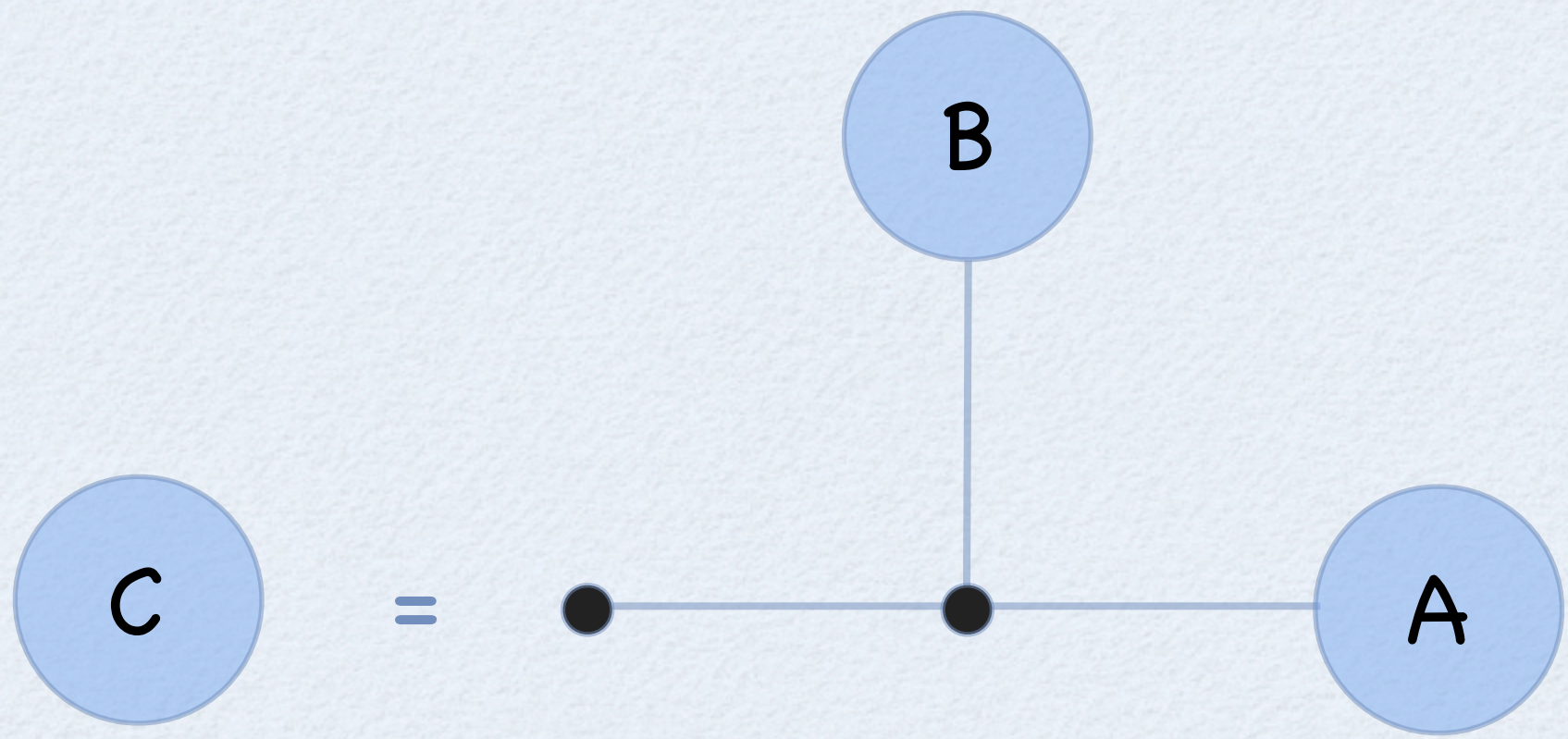
True both for expectation values and a.s. for a single comb. The right scaling limit is

$$\langle \tilde{Q}(\xi, \lambda) \rangle = \lim_{a \rightarrow 0} a^{1/2} \langle Q(a\xi, (\lambda/a)^{1-\alpha/2}) \rangle$$

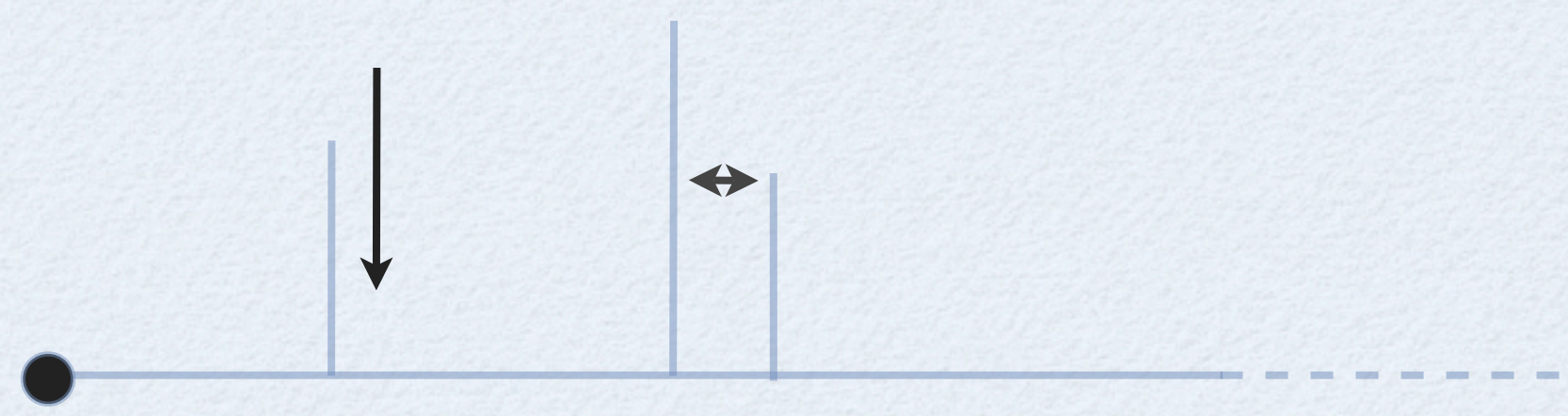
First return probabilities
related by

$$P_C = \frac{1-x}{3-P_A-P_B}$$

$$Q_G(x) = \frac{1}{1-P_G(x)}$$



Jensen's inequality gives lower bound



Graph moves give upper bound

Altogether we get

$$\frac{c}{\xi^{1/2} (1 + b (\lambda \xi)^{\alpha/2-1})^{1/2}} < Q(\xi, \lambda) < \frac{1}{\xi^{1/2}} F(\lambda \xi)$$

with

$$F(v) \sim \begin{cases} \text{const}, & v \rightarrow \infty \\ v^{\alpha/4}, & v \rightarrow 0 \end{cases}$$

$$\xi \rightarrow 0, d_s = 2 - \alpha/2$$

$$\xi \rightarrow \infty, d_s = 1$$

6. CONCLUSIONS

1. Spectral dimension can be defined on different scales
2. Order of limits is very important
3. We have considered only “kinematics” -- the probabilities $\pi(n)$ were put in by hand
4. Interesting to find solvable models where the affect appears dynamically